



# Li<sub>5</sub>Ti<sub>2</sub>O<sub>6</sub>F: a new low-loss oxyfluoride microwave dielectric ceramic for LTCC applications

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Received: 10 July 2019

Accepted: 11 September 2019

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## ABSTRACT

Use of microwave dielectric ceramics can significantly promote the development of communication devices. The primary criteria for determining their applications are excellent microwave dielectric properties and suitability for co-firing with cheap metals (e.g., Ag, Al). This study first reports the low-temperature synthesis of Li<sub>5</sub>Ti<sub>2</sub>O<sub>6</sub>F with a cubic rock salt structure. Li<sub>5</sub>Ti<sub>2</sub>O<sub>6</sub>F was prepared at 880 °C and had  $\epsilon_r$  of 19.6,  $Q \times f$  of 79500 GHz, and  $\tau_f$  of  $-29.6$  ppm/°C. The infrared reflectivity spectrum revealed that the dielectric contribution of Li<sub>5</sub>Ti<sub>2</sub>O<sub>6</sub>F in the microwave range was mainly affected by phonon absorption. Experimental study indicated that Li<sub>5</sub>Ti<sub>2</sub>O<sub>6</sub>F had good chemical compatibility versus Ag electrode at 880 °C for 2 h. Given these characteristics, Li<sub>5</sub>Ti<sub>2</sub>O<sub>6</sub>F ceramics exhibited potential for low-temperature co-fired ceramic technology.

## Introduction

Low-temperature co-fired ceramic (LTCC) is an essential material for filters, oscillators, and antenna substrates, due to its fast transmission speed, powerful function, small size, and lightweight [1–6]. The requirements to obtain LTCCs with crucial performance are as follows: (1) optimal sintering temperature below the melting point of the inner metal electrode such as Ag (961 °C) and good chemical compatibility with the metal; (2) suitable range of

relative permittivity ( $\epsilon_r$ ); (3) high quality factor ( $Q \times f$ ); (4) close-to-zero temperature coefficient of resonant frequency ( $\tau_f$ ); (5) low cost of raw materials; and (6) low bulk density for lightness of the electronic devices [7–13].

Commercial microwave dielectric ceramics are made of high-temperature oxides, such as CaTiO<sub>3</sub>, Al<sub>2</sub>O<sub>3</sub>, Ba(Mg<sub>1/3</sub>Ta<sub>2/3</sub>)O<sub>3</sub>, and Ba(Zn<sub>1/3</sub>Ta<sub>2/3</sub>)O<sub>3</sub>, which have sintering temperatures of about 1300 °C or even higher [14–17]. In addition, rock salt ceramics (Li<sub>2</sub>TiO<sub>3</sub>, Li<sub>2</sub>Mg<sub>2</sub>TiO<sub>5</sub>, and Li<sub>2</sub>Mg<sub>3</sub>SnO<sub>6</sub>) have good

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microwave dielectric properties and relatively high sintering temperatures ( $> 1200\text{ }^{\circ}\text{C}$ ) [18–20]. The most common method for decreasing the sintering temperature of ceramics is to add low-melting glasses, such as  $\text{Li}_2\text{O}$ – $\text{MgO}$ – $\text{B}_2\text{O}_3$ ,  $\text{B}_2\text{O}_3$ – $\text{Bi}_2\text{O}_3$ – $\text{SiO}_2$ – $\text{ZnO}$ , and  $\text{BaCu}(\text{B}_2\text{O}_5)$  [21–24]. However, this process is accompanied by reduced microwave dielectric properties. In recent years, fluoride microwave dielectric ceramics with low sintering temperatures have attracted research attention. Lei et al. [25] reported various low-fired fluoride microwave dielectric ceramics ( $\text{LiF}$ ,  $\text{CaF}_2$ ,  $\text{SrF}_2$ , and  $\text{BaF}_2$ ); in particular,  $\text{LiF}$  exhibits satisfactory performance ( $\epsilon_r = 8.02$ ,  $Q \times f = 73880\text{ GHz}$ ,  $\tau_f = -117.7\text{ ppm}/^{\circ}\text{C}$ ) at  $800\text{ }^{\circ}\text{C}$  and thus could be used to develop LTCC. In this study, a new type of oxyfluoride ceramic, namely  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$ , which possesses the advantages of low sintering temperature and high quality factor, was fabricated. The crystal structure, sintering characteristic, microwave dielectric properties, and chemical compatibility of  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  with the Ag electrode were investigated.

## Experimental procedure

Oxyfluoride  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  samples were fabricated through a conventional solid-state reaction in the air by using high-purity raw materials including  $\text{Li}_2\text{CO}_3$  (99.99%),  $\text{TiO}_2$  (99.99%), and  $\text{LiF}$  (99.99%), which were weighed via the correct stoichiometry. The mixed powders were ball-milled in a nylon jar with  $\text{ZrO}_2$  balls for 6 h by using ethanol as a medium at 180 rpm. The slurries were dried and calcined at  $700\text{ }^{\circ}\text{C}$  for 2 h. The calcined powders were re-milled for 6 h, dried, added with 5 wt% PVA as binder, and granulated. The resultant particles were pressed into cylinders with a diameter of 10 mm and a height of 6 mm and were heated to  $550\text{ }^{\circ}\text{C}$  for 4 h at a rate of  $1.5\text{ }^{\circ}\text{C}/\text{min}$  to remove the organic binder. Heating was continued up to  $800\text{ }^{\circ}\text{C}$ – $900\text{ }^{\circ}\text{C}$  for sintering. Powders of the same composition were covered in pellets and heated in an alumina crucible to limit Li evaporation during sintering. The chemical compatibility of  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  and Ag was also investigated by co-firing the mixtures at  $880\text{ }^{\circ}\text{C}$  for 2 h.

The bulk densities of the sintered samples were measured by Archimedes method. Thermal behavior and phase evolution were evaluated using a NETZSCH STA449F3 thermogravimetric and

differential scanning calorimetry analyzer (TG–DSC). The phase composition of  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  was determined using an X-ray diffractometer ( $\text{CuK}\alpha_1$ ,  $1.54059\text{ \AA}$ , model X'Pert PRO, PANalytical, Almelo, The Netherlands). The refinement XRD scan rate is  $0.021154^{\circ}/\text{s}$ . Energy-dispersive x-ray spectroscopy (EDS) and surface microstructures analyses of the specimens were performed using a field-emission scanning electron microscope (FE-SEM, model S4800, Hitachi, Japan). Selected area electron diffraction (SAED) patterns and high-resolution transmission electron microscope (HRTEM) images were collected by JEOL JEM-2100F TEM. The  $\epsilon_r$  and  $Q \times f$  of the sintered samples were measured in the  $\text{TE}_{011}$  mode according to the Hakki–Coleman method [26] using a network analyzer (model N5230A, Agilent Co., Palo Alto, CA).  $\tau_f$  was measured using a temperature chamber (Delta 9039, Delta Design, San Diego, CA) and calculated using the following equation:

$$\tau_f = \frac{f_{85} - f_{25}}{(85 - 25) \times f_{25}} \times 10^6 \quad (1)$$

where  $f_{85}$  and  $f_{25}$  are the resonant frequencies at  $85\text{ }^{\circ}\text{C}$  and  $25\text{ }^{\circ}\text{C}$ , respectively.

Room-temperature infrared reflectivity spectra were recorded using a Bruker IFS 66v FTIR spectrometer on infrared beamline station (U4) at National Synchrotron Radiation Lab (NSRL), Hefei, China. The infrared reflectivity spectra were analyzed by the harmonic oscillator model (the Lorentz three-parameter classical model) as follows: [27, 37]

$$\epsilon^*(\omega) = \epsilon_{\infty} + \sum_{j=1}^n \frac{\omega_{pj}^2}{\omega_{oj}^2 - \omega^2 - j\omega\gamma_j} \quad (2)$$

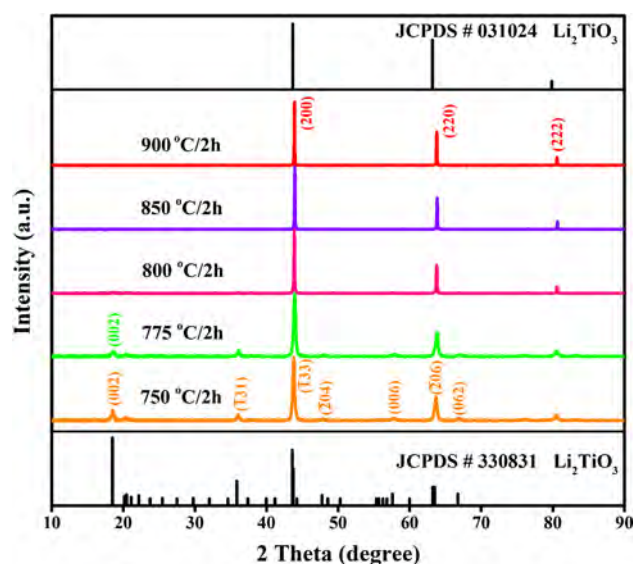
where  $\epsilon^*(\omega)$  is the complex dielectric function,  $\epsilon_{\infty}$  is the dielectric constant caused by the electronic polarization at high frequencies,  $\gamma_j$ ,  $\omega_{oj}$ , and  $\omega_{pj}$  are the damping factor, the transverse frequency, and plasma frequency of the  $j$ th Lorentz oscillator, respectively, and  $n$  is the number of transverse phonon modes. The complex reflectivity  $R(\omega)$  can be written as:

$$R(\omega) = \left| \frac{1 - \sqrt{\epsilon^*(\omega)}}{1 + \sqrt{\epsilon^*(\omega)}} \right|^2 \quad (3)$$

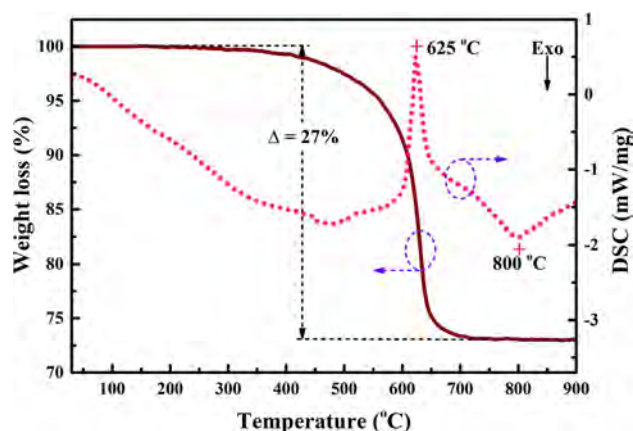
## Results and discussion

Figure 1 presents the room-temperature XRD patterns of  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  powders calcined at various temperatures for 2 h. The diffraction peak was indexed to the monoclinic  $\text{Li}_2\text{TiO}_3$  phase (JCPDS#330831) with the C2/c space group when calcined at 750 °C. The intensity of the (002) supercell diffraction peak decreased with increasing calcination temperature and basically disappeared when calcined at 800 °C. Hence, phase evolution occurred from order monoclinic to disorder cubic structure. This phenomenon confirmed with the results of the TG–DSC analysis (Fig. 2). At temperatures above 800 °C, the peaks were consistent with the cubic  $\text{Li}_2\text{TiO}_3$  phase (JCPDS#031024, Fm-3m space group). No trace of an impurity phase was detected.

Figure 2 shows the TG–DSC curves of the dried mixed raw material powders from 30 to 900 °C. In the DSC curve, an endothermic peak existed between 500 and 650 °C, corresponding to a mass loss of about 27% on the TG curve due to the thermal decomposition of  $\text{Li}_2\text{CO}_3$ . According to the phase diagram of  $\text{Li}_2\text{CO}_3$ –LiF [28], when  $\text{Li}_2\text{CO}_3$  and LiF were mixed in 2:1 mol ratio,  $\text{Li}_2\text{CO}_3$  was decomposed at 630 °C, consistent with our test result (625 °C). The exothermic peak at 800 °C was attributed to the phase transition of the monoclinic to the cubic structure and was verified from the XRD patterns.



**Figure 1** XRD patterns of  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  powders calcined at various temperatures.



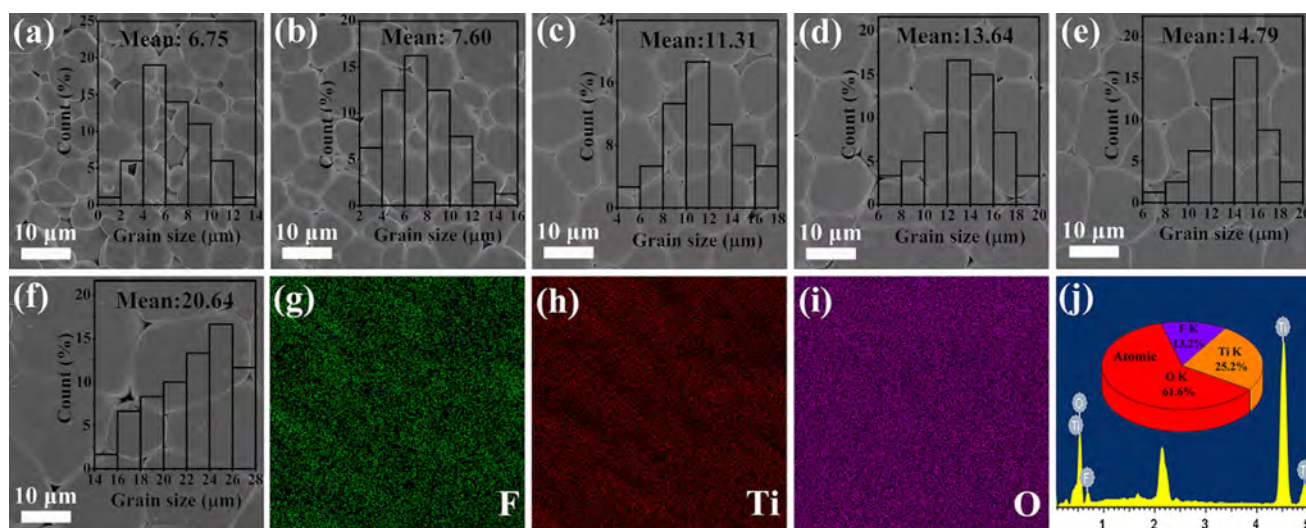
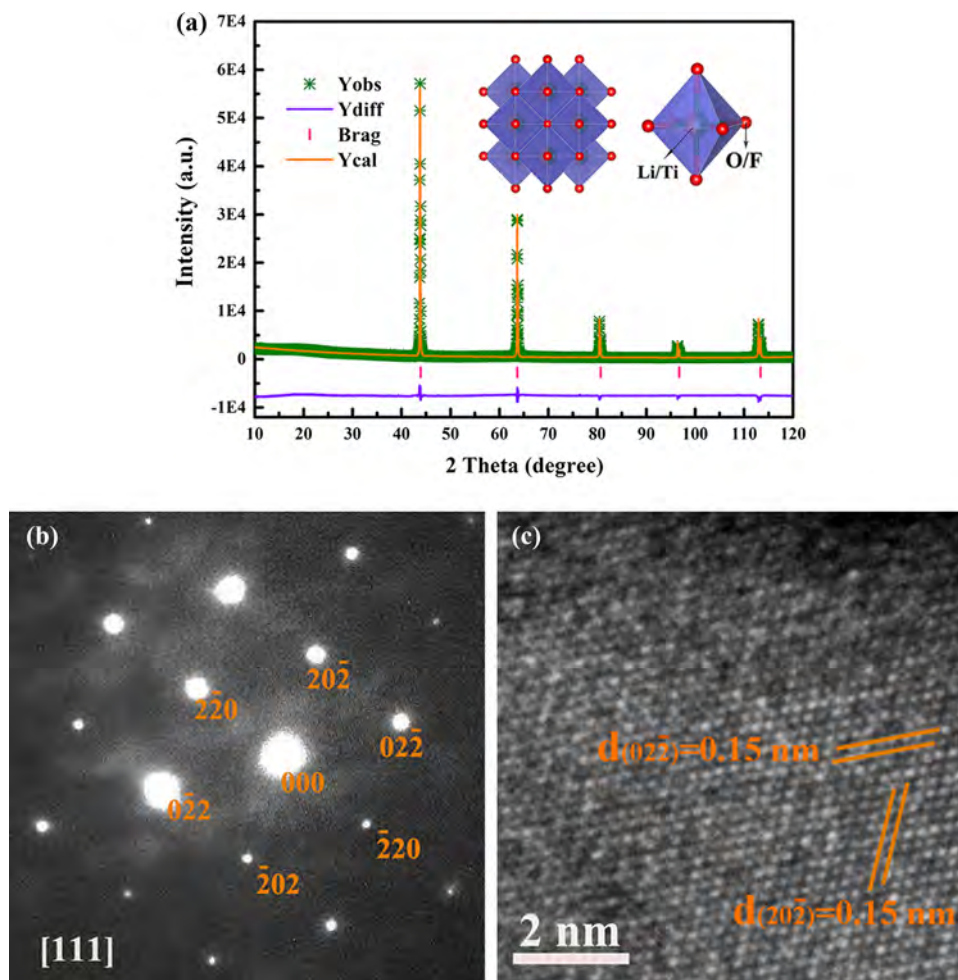
**Figure 2** TG–DSC curves of  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  powders from 30 to 900 °C.

Rietveld refinement of XRD was carried out using FullProf software to study the crystal structure of  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$ , as shown in Fig. 3a. The refinement confirmed that  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  had a cubic structure with a space group Fm-3m. The results of the refined parameters were  $a = b = c = 4.1325(1)$  Å,  $\alpha = \beta = \gamma = 90^\circ$ , and  $V_m = 70.5730(5)$  Å<sup>3</sup>, whereas the refinement reliability factors  $R_p$ ,  $R_{wp}$ , and  $R_{exp}$  were 4.78%, 8.44%, and 6.2%, respectively, indicating a reliable result. The schematic of the crystal structure of rock-salt-type  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  is illustrated in Fig. 3a. The O/F anions randomly stacking in a cubic close packing with Li/Ti cations occupied the octahedron sites. SAED patterns and the corresponding HRTEM images were recorded along the [111] zone axis to further verify the structure of  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  ceramic (Fig. 3b, c). The lattice fringes of the sample were 0.15 and 0.15 nm for the  $d_{(022)}$  and  $d_{(202)}$  spacing, respectively, which corresponded with the cubic rock salt structure.

Figure 4a–f shows the surface microstructures of  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  ceramics sintered at various temperatures. The grain size increased, whereas the size of pores decreased with increasing temperature. Moreover, a dense and homogeneous microstructure was obtained in the samples sintered at 880 °C (Fig. 4e). The grain size distributions and the average grain size of  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  are given in the inset of Fig. 4a–f. The average grain size increased from 6.75 to 14.79 μm when the temperature increased from 800 to 880 °C. However, abnormally large grains and a small number of pores were observed in the sample sintered at 900 °C. Element mapping of the samples sintered at 880 °C was conducted to confirm the



**Figure 3** **a** Rietveld refinement of samples sintered at 880 °C and schematic of the crystal structure for  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$ . **b** SAED pattern of  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  sample sintered at 880 °C. **c** HRTEM images of  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  ceramic recorded along the  $[111]$  zone axis.



**Figure 4** FE-SEM images of the surfaces of the sintered ceramic at: **a** 800 °C, **b** 820 °C, **c** 840 °C, **d** 860 °C, **e** 880 °C, and **f** 900 °C for  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  and corresponding EDS analysis surface scanning of **g** F, **h** Ti, **i** O, **j** EDS spectrum.

element distribution (Fig. 4g–i). All the elements, namely F, Ti, and O, were distributed

homogeneously. Element enrichment was not observed at the grain boundary, indicating that LiF

with  $\text{Li}_2\text{TiO}_3$  completely formed a solid solution. In addition, EDS analysis of the surface of the  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  ceramics sintered at 880 °C was also performed (Fig. 4j). The energy spectrum confirmed the existence of Ti, O, and F, and the ratio of each element was relatively close to the nominal composition.

The variations in the bulk density and relative density of the  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  ceramic sintered at 800 °C–900 °C are depicted in Fig. 5a. With increasing sintering temperature, the bulk density gradually increased and reached the saturated value of 3.12 g/cm<sup>3</sup> at 880 °C. The relative density of 94% slightly decreased with increasing sintering temperature. Figure 5b presents the  $\epsilon_r$  of  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  ceramic as a function of sintering temperature. The variation tendency of  $\epsilon_r$  was similar to that of relative density, that is, the influence of porosity on  $\epsilon_r$  plays an important role.  $\epsilon_r$  can then be corrected by the Bosman–Havenga formula [29]:

$$\epsilon_{\text{corr}} = \epsilon_r(1 + 1.5P) \quad (4)$$

$$P = 1 - \frac{\rho_{\text{mea}}}{\rho_{\text{th}}} \quad (5)$$

where  $P$  denotes the fractional porosity;  $\rho_{\text{mea}}$  and  $\rho_{\text{th}}$  are the measured and theoretical values of density,

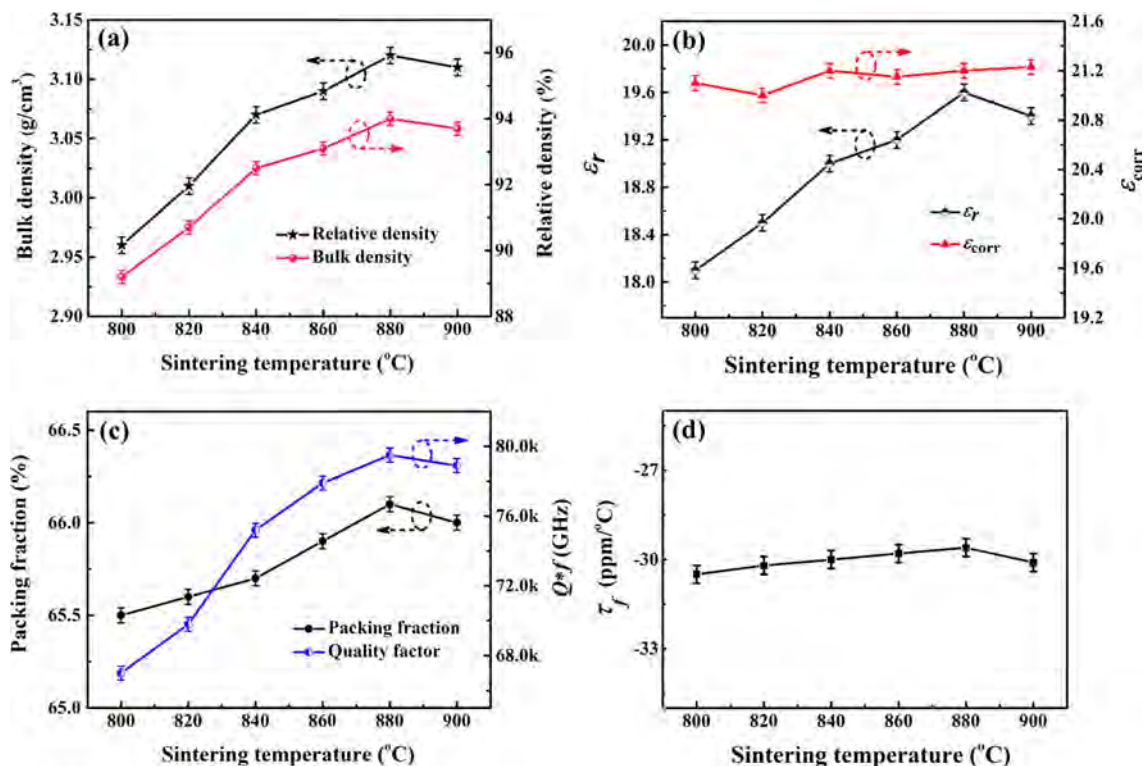
respectively. The corrected permittivity ( $\epsilon_{\text{corr}}$ ) values (Fig. 5b) fluctuate at around 21, which was higher than the measured value (18.1–19.6) of the samples sintered at 800 °C–900 °C. Furthermore, the theoretical permittivity ( $\epsilon_{\text{th}}$ ) can be computed using total ionic polarizability of individual ions and the molar volume of a compound in accordance with the Clausius–Mossotti equation: [30]

$$\epsilon_{\text{th}} = \frac{3V_m + 8\pi\alpha_D^T}{3V_m - 4\pi\alpha_D^T} \quad (6)$$

where  $V_m$  and  $\alpha_D^T$  represent cell volume and molecular polarizability, respectively. The sum of ionic polarizability of  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  can be calculated with the additive rule:

$$\alpha(\text{Li}_5\text{Ti}_2\text{O}_6\text{F}) = \frac{5}{7}\alpha(\text{Li}^+) + \frac{2}{7}\alpha(\text{Ti}^{4+}) + \frac{6}{7}\alpha(\text{O}^{2-}) + \frac{1}{7}\alpha(\text{F}^-) \quad (7)$$

where  $\alpha(\text{Li}^+)$ ,  $\alpha(\text{Ti}^{4+})$ ,  $\alpha(\text{O}^{2-})$ , and  $\alpha(\text{F}^-)$  are 1.2, 2.93, 2.01, and 1.62 Å<sup>3</sup>, respectively [29, 31]. According to Eq. (6),  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  ceramic has a theoretical permittivity of 20.7, which is comparable to the measured value.



**Figure 5** a Bulk density and relative density, b relative permittivity ( $\epsilon_r$  and  $\epsilon_{\text{corr}}$ ), c packing fraction and quality factor ( $Q \times f$ ), and d temperature coefficient of resonance frequency ( $\tau_f$ ) of  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  ceramics sintered at 800–900 °C for 2 h.

The  $Q \times f$  values of  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  ceramic increased from 67000 to 79500 GHz as the sintering temperature was increased from 800 to 880 °C (Fig. 5c). In general,  $Q \times f$  value was influenced by external factors, such as density, impurity, and secondary phase, and internal factors, such as crystal structure [32, 33]. Packing fraction theory was proposed to explain the quality factor and defined as Eq. (8) [34]:

$$\text{packing fraction(\%)} = \frac{\text{volume of packed ions}}{\text{volume of unit cell}} \times Z = \frac{\left(\frac{4\pi}{3}\right)(r_{\text{Li}}^3 + r_{\text{Ti}}^3 + r_{\text{O}}^3 + r_{\text{F}}^3)}{V} \times 4 \quad (8)$$

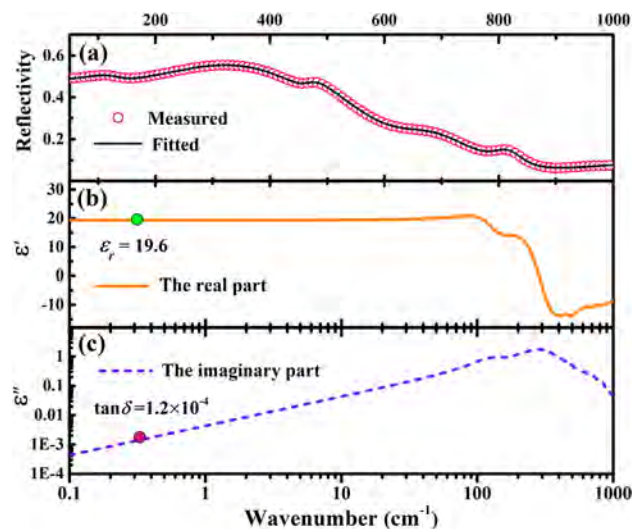
The packing fraction was positively correlated with  $Q \times f$ , as shown in Fig. 5c. An increase in the packing fraction reduced the vibration of the crystal lattice, causing an increase in the quality factor [35]. The maximum calculated packing fraction was 66.1% when  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  ceramic was sintered at 880 °C. As shown in Fig. 5d, the  $\tau_f$  values of  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  ceramic were not temperature dependent and remained stable at around  $-30$  ppm/°C. In summary,  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  ceramic sintered at 880 °C exhibited the optimum microwave dielectric properties with  $\epsilon_r = 19.6$ ,  $Q \times f = 79500$  GHz, and  $\tau_f = -29.6$  ppm/°C.

Infrared reflectivity spectrum was used to investigate inherent dielectric properties [36, 37]. Figure 6 shows the IR reflectivity spectrum of  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  ceramic sintered at 880 °C. The fitting of reflectivity spectrum was performed by REFFIT software, and processing procedure is similar as described in our previous work [12, 36]. As listed in Table 1, the relative permittivity at high frequencies ( $\epsilon_\infty$ ) was 5.71, and the intrinsic permittivity  $\epsilon_0$  (extrapolated value) was 19.4, which was close to the measured value of 19.6. The permittivity of  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  ceramic was mainly attributed to the dielectric response of intrinsic phonons. Furthermore, dielectric loss can be calculated by the following formula:

$$\epsilon' = \epsilon_\infty + \sum_{j=1}^n \frac{\omega_{pj}^2}{\omega_{oj}^2} = \epsilon_\infty + \sum_{j=1}^n \Delta\epsilon_j \quad (9)$$

$$\tan \delta = \frac{\epsilon''}{\epsilon'} = \omega \sum_{j=1}^n \frac{\Delta\epsilon_j \gamma_j}{\omega_{oj}^2 (\epsilon_\infty + \sum_{j=1}^n \Delta\epsilon_j)} \quad (10)$$

The calculated dielectric loss was  $9.8 \times 10^{-5}$  ( $f = 9.5$  GHz,  $Q = 1/\tan \delta$ ,  $Q \times f = 97000$  GHz), which was lower than the measured value of



**Figure 6** a Measured and fitted IR reflectivity spectra of the  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  ceramic within 50–1000  $\text{cm}^{-1}$ . b The yellow solid line is the real part. c The purple dotted line is the imaginary part.

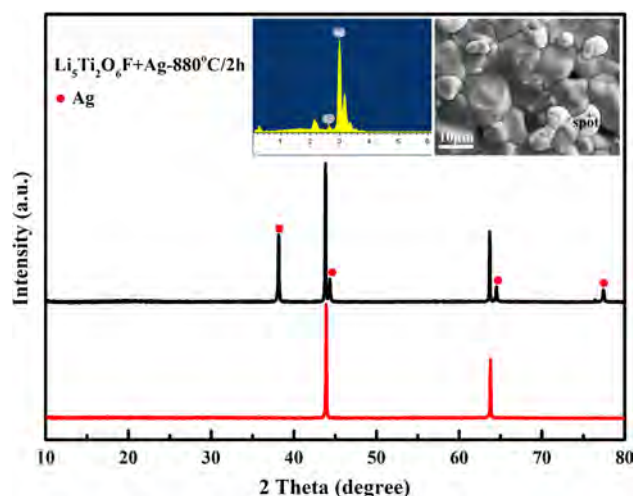
**Table 1** Phonon parameters obtained by fitting of the infrared reflectivity spectrum of  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  ceramic

Mode	$\omega_{oj}$	$\omega_{pj}$	$\gamma_j$	$\Delta\epsilon_j$
1	128.18	252.8	75.605	1.6740
2	306.58	1376.6	234.66	17.200
3	470.29	215.3	55.175	0.2100
4	489.7	323.22	276.43	0.4360
5	677.58	501.51	194.50	0.5480
6	812.56	226.45	63.506	0.0747
$\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$ $\epsilon_\infty = 5.71$ $\epsilon_0 = 19.4$				

$1.2 \times 10^{-4}$  ( $Q \times f = 79500$  GHz). The deviation between the calculated and measured quality factor might be attributed to extrinsic factors, such as pores and uneven distribution of grain size of the sintered sample. Therefore, quality factor can be improved by reducing the porosity and optimizing the preparation process.

$\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  powders were co-fired with Ag powders to study the chemical compatibility of ceramics with metal electrodes. Figure 7 shows the XRD pattern, EDS, and BSEM images of pure  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  and  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  co-fired with 20 wt% Ag powders sintered at 880 °C for 2 h. The peaks of  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  and Ag phases were observed, indicating their good chemical compatibility. Two kinds of grains with different sizes and shapes were detected in the BSE image. The brighter grains marked with “spot” were identified





**Figure 7** XRD and BSEM analyses of  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  ceramic co-fired with 20 wt% Ag at 880 °C for 2 h.

through elemental analysis with EDS as Ag. All the results indicated the lack of chemical reaction between  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  and Ag; as such,  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  exhibits potential for application in LTCC technology.

## Conclusions

In this work, a new oxyfluoride,  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$ , was synthesized using traditional solid-state reaction method at 800 °C. The XRD Rietveld refinement revealed that  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  belongs to a cubic structure with a space group  $\text{Fm-3m}$ . The SAED patterns and corresponding HRTEM images confirmed the structure. After sintering at 880 °C,  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  ceramic exhibited the optimum microwave dielectric properties:  $\epsilon_r = 19.6$ ,  $Q \times f = 79500 \text{ GHz}$ , and  $\tau_f = -29.6 \text{ ppm/}^\circ\text{C}$ . The results of the far-infrared reflectivity test revealed that the polarization of the  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  ceramic in the microwave region was mainly attributed to the absorption of phonon oscillation in the infrared region and that the defect phonon scattering had minimal effect. Moreover,  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  ceramic was found to be chemically compatible with Ag. Hence,  $\text{Li}_5\text{Ti}_2\text{O}_6\text{F}$  ceramic is a potential candidate for LTCC applications.

## Acknowledgements

We appreciate the administrators of the IR beamline workstation of the National Synchrotron Radiation

Laboratory (NSRL) for their help in IR measurements. This work was supported by the Natural Science Foundation of China (Nos. 21561008, 21965009), the Natural Science Foundation of Guangxi Zhuang Autonomous Region (Nos. 2015GXNSFFA139003, 2018GXNSFAA138175), Project of Scientific Research and Technical Exploitation Program of the Guangxi Zhuang Autonomous Region (Nos. AA18118008, AA18118034, AA18118023) and Guilin (20170225), and Innovation Project of Guangxi Graduate Education (YCSW2019157).

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Dalton Trans 42:1483–1491

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